

**Example 136.** Spell out the Taylor method of order 3 for numerically solving the IVP

$$y' = \cos(x)y, \quad y(0) = 1.$$

**Solution.** The Taylor method of order 3 is based on the Taylor expansion

$$y(x+h) = y(x) + y'(x)h + \frac{1}{2}y''(x)h^2 + \frac{1}{6}y'''(x)h^3 + O(h^4),$$

where we have a local truncation error of  $O(h^4)$  so that the global error will be  $O(h^3)$ .

From the DE we know that  $y'(x) = \cos(x)y$ , which is  $f(x, y)$ . As in the previous example, we differentiate this to obtain

$$\begin{aligned} y''(x) &= \frac{d}{dx} \cos(x)y = -\sin(x)y + \cos(x)y' = -\sin(x)y + \cos^2(x)y \\ &= (-\sin(x) + \cos^2(x))y, \end{aligned}$$

which is  $f'(x, y)$ . Once more differentiating this, we obtain

$$\begin{aligned} y'''(x) &= \frac{d}{dx} (-\sin(x) + \cos^2(x))y \\ &= (-\cos(x) - 2\cos(x)\sin(x))y + (-\sin(x) + \cos^2(x))y' \\ &= (\cos^2(x) - 3\sin(x) - 1)\cos(x)y, \end{aligned}$$

which is  $f''(x, y)$ . Hence, the Taylor method of order 3 takes the form:

$$\begin{aligned} y_{n+1} &= y_n + f(x_n, y_n)h + \frac{1}{2}f'(x_n, y_n)h^2 + \frac{1}{6}f''(x_n, y_n)h^3 \\ &= y_n + \cos(x_n)y_n h + \frac{1}{2}((-\sin(x_n) + \cos^2(x_n))y_n)h^2 + \frac{1}{6}((\cos^2(x_n) - 3\sin(x_n) - 1)\cos(x_n)y_n)h^3 \end{aligned}$$

**Comment.** The exact solution of this IVP is  $y(x) = e^{\sin(x)}$ . We use this in the next example for comparison.

**Example 137.** Python Let us use Python to approximate the solution  $y(x)$  of the IVP from the previous example for  $x \in [0, 2]$ .

```
>>> from math import e, cos, sin
>>> def taylor_3_cosy(x0, y0, xmax, n):
    h = (xmax - x0) / n
    ypoints = [y0]
    for i in range(n):
        y0 = y0 + cos(x0)*y0*h + 1/2*(cos(x0)**2-sin(x0))*y0*h**2 + \
            1/6*(cos(x0)**2-3*sin(x0)-1)*cos(x0)*y0*h**3
        x0 = x0 + h
        ypoints.append(y0)
    return ypoints
```

Since the exact solution is  $y(x) = e^{\sin(x)}$ , we have  $y(2) = e^{\sin(2)} \approx 2.48258$ .

```
>>> taylor_3_cosy(0, 1, 2, 4)
[1, 1.625, 2.3475297541746047, 2.7350418255304874, 2.476391322837691]
```

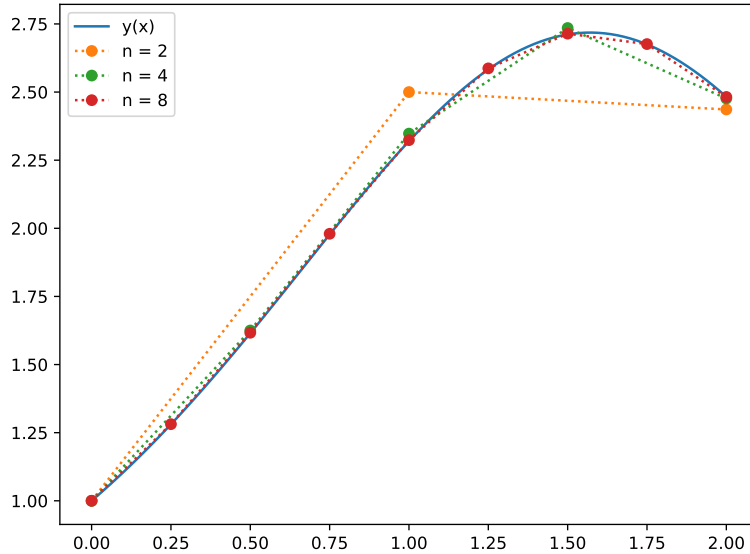
For comparison, the exact values of the solution at these four steps are:

```
>>> [e**sin(i/2) for i in range(5)]
[1.0, 1.6151462964420837, 2.319776824715853, 2.7114810176821584, 2.4825777280150003]
```

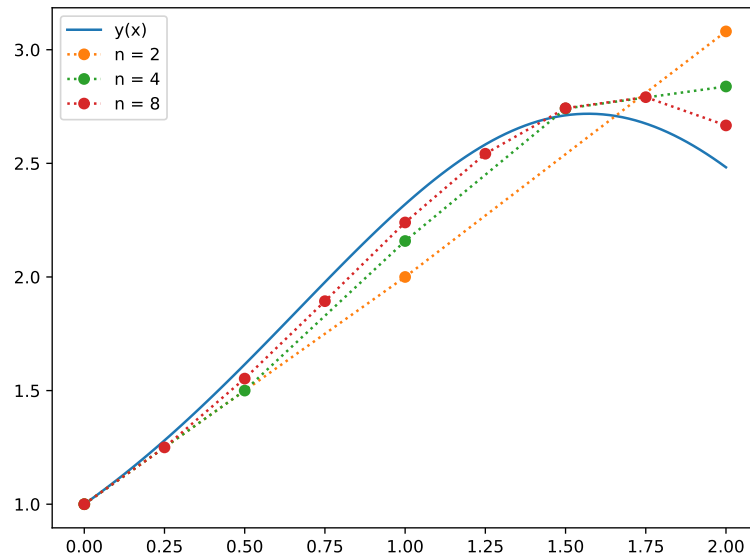
The following convincingly illustrates that the error is indeed  $O(h^3)$ .

```
>>> [taylor_3_cosy(0, 1, 2, 10**n)[-1] - e**sin(2) for n in range(5)]  
[2.5174222719849997, -0.0002461575553160955, -1.6375769584797695e-07, -  
1.5647971807197791e-10, 2.0339285811132868e-13]
```

The following is a plot of the exact solution together with our approximations when using 2, 4 and 8 steps. Already for 4 steps, we obtain an approximation that, at least visually, is remarkably good.



On the other hand, for comparison, the following plot shows the corresponding approximations when using Euler's method instead.



## Midpoint method (also known as the modified Euler method)

Recall that, in Euler's method (which is the same as the Taylor method of order 1) with step size  $h$ , we use  $y_{n+1} = y_n + f(x_n, y_n)h$  because  $f(x_n, y_n)$  is an approximation of  $y'(x_n)$ .

To get a better approximation  $\tilde{y}_{n+1}$ , we could do two small Euler steps (of size  $h/2$  each) instead. We can then use the idea of Richardson extrapolation to combine Euler's  $y_{n+1}$  and the twice-Euler  $\tilde{y}_{n+1}$  to get an approximation of higher order. This results in the following method:

**(midpoint method)** The following is an order 2 method for solving IVPs:

$$y_{n+1} = y_n + f\left(x_n + \frac{h}{2}, y_n + f(x_n, y_n)\frac{h}{2}\right)h$$

**Advantage over Taylor method.** When compared with the Taylor method of order 2, this has the advantage of not requiring us to determine partial derivatives of  $f(x, y)$ .

**Why "midpoint"?** In Euler's method, we use  $y'(x_n) \approx f(x_n, y_n)$  as the slope for stepping from  $x_n$  to  $x_{n+1} = x_n + h$ . That is a bit unbalanced since we use the slope at the left endpoint for the entire interval (that's why Euler's method in the form we have seen is often called the **forward Euler method**). Note that the midpoint method is more balanced because it uses (an approximation of) the derivative at the midpoint  $x_n + h/2$  instead.

**How to derive the midpoint method.** As mentioned above, we can derive the midpoint method by applying the idea of Richardson extrapolation to the Euler method.

- If we do one Euler step, then our approximation of  $y(x_{n+1})$  is  $y_{n+1}^{(1)} = y_n + f(x_n, y_n)h$ . We know that the error in this approximation is roughly  $Ch^2$  (in fact,  $C \approx \frac{1}{2}y''(x_n)$  if  $h$  is small).
- If we do two Euler steps, then our approximation is  $y_{n+1}^{(2)} = y_{\text{half}} + f(x_{\text{half}}, y_{\text{half}})\frac{h}{2}$  where  $x_{\text{half}} = x_n + \frac{h}{2}$  and  $y_{\text{half}} = y_n + f(x_n, y_n)\frac{h}{2}$  are the result of the first Euler step with step size  $\frac{h}{2}$ . Combined, we get  $y_{n+1}^{(2)} = y_n + f(x_n, y_n)\frac{h}{2} + f\left(x_n + \frac{h}{2}, y_n + f(x_n, y_n)\frac{h}{2}\right)\frac{h}{2}$ . The error should be roughly  $2C\left(\frac{h}{2}\right)^2 = \frac{1}{2}Ch^2$  because we are doing 2 steps with step size  $\frac{h}{2}$ . [For small  $h$ , the constant  $C$  should still be as before, at least to first order.]
- As for Richardson extrapolation,  $2y_{n+1}^{(2)} - y_{n+1}^{(1)}$  should therefore be an approximation of  $y(x_{n+1})$  of order 3 (for the local truncation error). Indeed,  $2y_{n+1}^{(2)} - y_{n+1}^{(1)} = y_n + f\left(x_n + \frac{h}{2}, y_n + f(x_n, y_n)\frac{h}{2}\right)h$  is precisely the midpoint method, which therefore should have global error of order 2.

**An alternative way to show that the order is 2.** We can also show that the midpoint method is of order 2 by computing and comparing Taylor expansions. The true solution has the expansion

$$\begin{aligned} y(x+h) &= y(x) + y'(x)h + \frac{1}{2}y''(x)h^2 + O(h^3) \\ &= y(x) + f(x, y(x))h + \frac{1}{2} \left[ \frac{d}{dx} f(x, y(x)) \right] h^2 + O(h^3) \\ &= y(x) + f(x, y)h + \frac{1}{2} (f_x(x, y) + f_y(x, y) \cdot y'(x)) h^2 + O(h^3). \end{aligned}$$

On the other hand, expanding the  $f(\dots)$  term on the right-hand side of the midpoint method (with  $x$  and  $y$  in place of  $x_n$  and  $y_n$ ) using the multivariate chain rule, we find

$$y + f\left(x + \frac{h}{2}, y + f(x, y)\frac{h}{2}\right)h = y + \left( f(x, y) + \left( f_x(x, y)\frac{1}{2} + f_y(x, y)\frac{f(x, y)}{2} \right)h + O(h^2) \right)h,$$

which agrees with the true expansion up to  $O(h^3)$ .

**Example 138.** Consider the IVP  $y' = y$ ,  $y(0) = 1$ . Approximate the solution  $y(x)$  for  $x \in [0, 1]$  using the midpoint method with 4 steps. In particular, what is the approximation for  $y(1)$ ?

**Comment.** Compare with Example 133. The real solution is  $y(x) = e^x$  so that  $y(1) = e \approx 2.71828$ .

**Solution.** The step size is  $h = \frac{1-0}{4} = \frac{1}{4}$ . We apply the midpoint method with  $f(x, y) = y$ :

$$\begin{aligned} x_0 = 0 & & y_0 = 1 \\ x_1 = \frac{1}{4} & & y_1 = y_0 + f\left(x_0 + \frac{h}{2}, y_0 + f(x_0, y_0)\frac{h}{2}\right)h = \frac{41}{32} \approx 1.2813 \\ x_2 = \frac{1}{2} & & y_2 = y_1 + f\left(x_1 + \frac{h}{2}, y_1 + f(x_1, y_1)\frac{h}{2}\right)h = \frac{41^2}{32^2} \approx 1.6416 \\ x_3 = \frac{3}{4} & & y_3 = y_2 + f\left(x_2 + \frac{h}{2}, y_2 + f(x_2, y_2)\frac{h}{2}\right)h = \frac{41^3}{32^3} \approx 2.1033 \\ x_4 = 1 & & y_4 = y_3 + f\left(x_3 + \frac{h}{2}, y_3 + f(x_3, y_3)\frac{h}{2}\right)h = \frac{41^4}{32^4} \approx 2.6949 \end{aligned}$$

In particular, the approximation for  $y(1)$  is  $y_4 \approx 2.6949$ , which is a noticeable improvement over Example 133, where we used the Euler method instead.

**Comment.** The above computations become particularly transparent if we realize that, for  $f(x, y) = y$ , each Euler step takes the simple form  $y_{n+1} = y_n + f\left(x_n + \frac{h}{2}, y_n + f(x_n, y_n)\frac{h}{2}\right)h = y_n\left(1 + h + \frac{h^2}{2}\right)$ .

It follows that  $y_n = \left(1 + h + \frac{h^2}{2}\right)^n$ . Can you see how, as  $h \rightarrow 0$ , this allows us to recover the exact solution?

**Example 139.** Python Let us apply the midpoint method to  $y' = y$ ,  $y(0) = 1$ .

```
>>> def midpoint(f, x0, y0, xmax, n):
    h = (xmax - x0) / n
    ypoints = [y0]
    for i in range(n):
        y0 = y0 + f(x0+h/2, y0+f(x0,y0)*h/2)*h
        x0 = x0 + h
        ypoints.append(y0)
    return ypoints

>>> def f_y(x, y):
    return y
```

The exact solution is  $y(x) = e^x$  with  $y(1) = e \approx 2.718$ .

```
>>> midpoint(f_y, 0, 1, 1, 4)

[1, 1.28125, 1.6416015625, 2.103302001953125, 2.6948556900024414]
```

The following numerically confirms that the error in the midpoint method is  $O(h^2)$ .

```
>>> from math import e
>>> [midpoint(f_y, 0, 1, 1, 10**n)[-1] - e for n in range(6)]

[-0.2182818284590451, -0.004200981850821073, -4.49658990882007e-05, -
4.5270728232793545e-07, -4.530157138304958e-09, -4.530020802917534e-11]
```

## Runge–Kutta methods

The midpoint method can be written as:

$$\begin{aligned}y_{n+1} &= y_n + K_1 h \\K_0 &= f(x_n, y_n) \\K_1 &= f\left(x_n + \frac{h}{2}, y_n + K_0 \frac{h}{2}\right)\end{aligned}$$

Note that replacing the rule by  $y_{n+1} = y_n + K_0 h$  results in Euler's method. Indeed, both  $K_0$  and  $K_1$  are approximations of the slope  $y'$  that we need for stepping from  $x_n$  to  $x_{n+1} = x_n + h$ .

Adding further such approximations  $K_i$  to the mix, one can eliminate further terms in the error expansion and obtain higher order methods known as **Runge–Kutta methods**.

The midpoint method is an example of a Runge–Kutta method of order 2 (but there are others as well).

[https://en.wikipedia.org/wiki/Runge%E2%80%93Kutta\\_methods](https://en.wikipedia.org/wiki/Runge%E2%80%93Kutta_methods)

Of particular practical importance is the following instance:

### (Runge–Kutta method of order 4)

$$\begin{aligned}y_{n+1} &= y_n + \frac{1}{6}(K_0 + 2K_1 + 2K_2 + K_3)h \\K_0 &= f(x_n, y_n) \\K_1 &= f\left(x_n + \frac{h}{2}, y_n + K_0 \frac{h}{2}\right) \\K_2 &= f\left(x_n + \frac{h}{2}, y_n + K_1 \frac{h}{2}\right) \\K_3 &= f(x_n + h, y_n + K_2 h)\end{aligned}$$

**Comment.** Note how each of  $K_0, K_1, K_2, K_3$  is an approximation of  $y'$  on the interval  $[x_n, x_{n+1}]$  (with  $K_0$  closer to  $y'(x_n)$  and  $K_3$  closer to  $y'(x_{n+1})$ ). By taking the appropriate weighted average, we are able to get an approximation with a higher order.

**Advanced comment.** Note that the weights (with  $K_1$  and  $K_2$  combined because they both correspond to the midpoint  $x_n + h/2$ ) are the same as in Simpson's rule for numerical integration. That is more than a coincidence. Indeed, if  $f(x, y) = f(x)$  does not depend on  $y$ , then solving the DE is equivalent to integrating  $f(x)$  and the Runge–Kutta method of order 4 turns into Simpson's rule.

**Example 140.** Python Let us implement the Runge–Kutta method of order 4.

```
>>> def runge_kutta4(f, x0, y0, xmax, n):
    h = (xmax - x0) / n
    ypoints = [y0]
    for i in range(n):
        K0 = f(x0, y0)
        K1 = f(x0+h/2, y0+K0*h/2)
        K2 = f(x0+h/2, y0+K1*h/2)
        K3 = f(x0+h, y0+K2*h)
        y0 = y0 + (K0 + 2*K1 + 2*K2 + K3)*h/6
        x0 = x0 + h
        ypoints.append(y0)
    return ypoints
```

First, for comparison with earlier methods, let us apply the method to the IVP  $y' = y$ ,  $y(0) = 1$ , which has the exact solution  $y(x) = e^x$  with  $y(1) = e \approx 2.718$ .

```
>>> def f_y(x, y):
    return y

>>> runge_kutta4(f_y, 0, 1, 1, 4)

[1, 1.2840169270833333, 1.648699469036526, 2.1169580259162033, 2.718209939201323]
```

The following convincingly illustrates that the error is indeed  $O(h^4)$ .

```
>>> from math import e

>>> [runge_kutta4(f_y, 0, 1, 1, 10**n)[-1] - e for n in range(6)]

[-0.009948495125712054, -2.0843238792700447e-06, -2.2464119453502462e-10, -
2.042810365310288e-14, 1.1546319456101628e-14, 6.217248937900877e-15]
```

Pause for a moment to really appreciate how much better these errors are in comparison with Euler's method! Whereas computing  $10^5$  values with Euler's method resulted in an error of  $1.36 \cdot 10^{-5}$ , we are now able to obtain an error of  $2.04 \cdot 10^{-14}$  with only  $10^3$  values.

As a second example, let us consider as in Example 137 the IVP  $y' = \cos(x)y$  with  $y(0) = 1$ , which has the exact solution  $y(x) = e^{\sin(x)}$  with  $y(2) = e^{\sin(2)} \approx 2.48258$ .

```
>>> def f_cosx_y(x, y):
    return cos(x)*y

>>> runge_kutta4(f_cosx_y, 0, 1, 2, 4)

[1, 1.614859377441316, 2.3191895982789603, 2.7107641474177457, 2.481902218021582]
```

The following again convincingly illustrates that the error is indeed  $O(h^4)$ .

```
>>> from math import e

>>> [runge_kutta4(f_cosx_y, 0, 1, 2, 10**n)[-1] - e**sin(2) for n in range(5)]

[-0.12999578105593113, -1.726387102785054e-05, -1.6494263732624859e-09, -
1.6431300764452317e-13, 3.419486915845482e-13]
```

**Important comment.** Note that, in contrast to Example 137, we did not have to compute partial derivatives of  $f(x, y) = \cos(x)y$  by hand. Instead, we were able to simply use  $\cos(x)y$  in our `runge_kutta4` function.