

**Example 131.** Find all eigenfunctions and eigenvalues of

$$y'' + \lambda y = 0, \quad y'(0) = 0, \quad y(3) = 0.$$

**Solution.** We distinguish three cases:

$\lambda < 0$ . The characteristic roots are  $\pm r = \pm\sqrt{-\lambda}$  and the general solution to the DE is  $y(x) = Ae^{rx} + Be^{-rx}$ . Then  $y'(0) = Ar - Br = 0$  implies  $B = A$ , so that  $y(3) = A(e^{3r} + e^{-3r})$ . Since  $e^{3r} + e^{-3r} > 0$ , we see that  $y(3) = 0$  only if  $A = 0$ . So there is no solution for  $\lambda < 0$ .

$\lambda = 0$ . The general solution to the DE is  $y(x) = A + Bx$ . Then  $y'(0) = 0$  implies  $B = 0$ , and it follows from  $y(3) = A = 0$  that  $\lambda = 0$  is not an eigenvalue.

$\lambda > 0$ . The characteristic roots are  $\pm i\sqrt{\lambda}$ . So, with  $r = \sqrt{\lambda}$ , the general solution is  $y(x) = A \cos(rx) + B \sin(rx)$ .  $y'(0) = Br = 0$  implies  $B = 0$ . Then  $y(3) = A \cos(3r) = 0$ . Note that  $\cos(3r) = 0$  is true if and only if  $3r = \frac{\pi}{2} + n\pi = \frac{(2n+1)\pi}{2}$  for some integer  $n$ . Since  $r > 0$ , we have  $n \geq 0$ . Correspondingly,  $\lambda = r^2 = \left(\frac{(2n+1)\pi}{6}\right)^2$  and  $y(x) = \cos\left(\frac{(2n+1)\pi}{6}x\right)$ .

In summary, we have that the eigenvalues are  $\lambda = \left(\frac{(2n+1)\pi}{6}\right)^2$ , with  $n = 0, 1, 2, \dots$  with corresponding eigenfunctions  $y(x) = \cos\left(\frac{(2n+1)\pi}{6}x\right)$ .

## Partial differential equations

### The heat equation

We wish to describe one-dimensional heat flow.

**Comment.** If this sounds very specialized, it might help to know that the heat equation is also used, for instance, in probability (Brownian motion), financial math (Black-Scholes), or chemical processes (diffusion equation).

Let  $u(x, t)$  describe the temperature at time  $t$  at position  $x$ .

If we model a heated rod of length  $L$ , then  $x \in [0, L]$ .

**Notation.**  $u(x, t)$  depends on two variables. When taking derivatives, we will use the notations  $u_t = \frac{\partial}{\partial t}u$  and  $u_{xx} = \frac{\partial^2}{\partial x^2}u$  for first and higher derivatives.

Experience tells us that heat flows from warmer to cooler areas and has an averaging effect.

Make a sketch of some temperature profile  $u(x, t)$  for fixed  $t$ .

As  $t$  increases, we expect maxima (where  $u_{xx} < 0$ ) of that profile to flatten out (which means that  $u_t < 0$ ); similarly, minima (where  $u_{xx} > 0$ ) should go up (meaning that  $u_t > 0$ ). The simplest relationship between  $u_t$  and  $u_{xx}$  which conforms with our expectation is  $u_t = ku_{xx}$ , with  $k > 0$ .

**(heat equation)**

$$u_t = ku_{xx}$$

Note that the heat equation is a linear and homogeneous **partial differential equation**.

In particular, the principle of superposition holds: if  $u_1$  and  $u_2$  solve the heat equation, then so does  $c_1u_1 + c_2u_2$ .

**Higher dimensions.** In higher dimensions, the heat equation takes the form  $u_t = k(u_{xx} + u_{yy})$  or  $u_t = k(u_{xx} + u_{yy} + u_{zz})$ . Note that  $\Delta u = u_{xx} + u_{yy} + u_{zz}$  is the Laplace operator you may know from Calculus III.

The Laplacian  $\Delta u$  is also often written as  $\Delta u = \nabla^2 u$ . The operator  $\nabla = (\partial/\partial x, \partial/\partial y)$  is pronounced "nabla" (Greek for a certain harp) or "del" (Persian for heart), and  $\nabla^2$  is short for the inner product  $\nabla \cdot \nabla$ .

**Example 132.** Note that  $u(x, t) = ax + b$  solves the heat equation.

**Example 133.** To get a feeling, let us find some other solutions to  $u_t = u_{xx}$  (for starters,  $k = 1$ ).

- For instance,  $u(x, t) = e^t e^x$  is a solution.  
[Not a very interesting one for modeling heat flow because it increases exponentially in time.]
- ... to be continued ...  
Can you find further solutions?

Let us think about what is needed to describe a unique solution of the heat equation.

- **Initial condition** at  $t = 0$ :  $u(x, 0) = f(x)$  (IC)

This specifies an initial temperature distribution at time  $t = 0$ .

- **Boundary condition** at  $x = 0$  and  $x = L$ : (BC)

Assuming that heat only enters/exits at the boundary (think of our rod as being insulated, except possibly at the two ends), we need some condition on the temperature at the ends. For instance:

- $u(0, t) = A, u(L, t) = B$

This models a rod where one end is kept at temperature  $A$  and the other end at temperature  $B$ .

- $u_x(0, t) = u_x(L, t) = 0$

This models a rod whose ends are insulated as well.

Under such assumptions, our physical intuition suggests that there should be a unique solution.

**Important comment.** We can always transform the case  $u(0, t) = A, u(L, t) = B$  into  $u(0, t) = u(L, t) = 0$  by using the fact that  $u(t, x) = ax + b$  solves  $u_t = k u_{xx}$ . Can you spell this out?